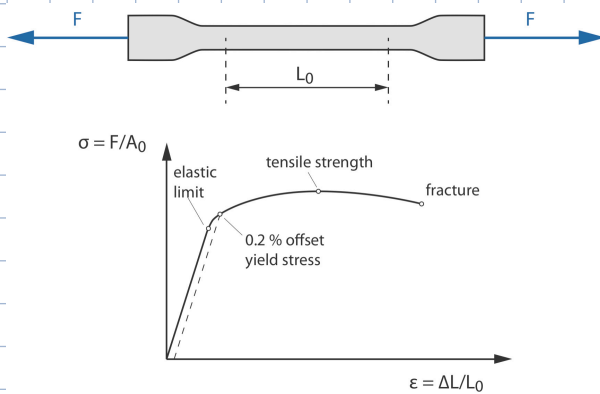


Governing equations of elasticity :

- equilibrium equations : stress transformation
- compatibility equations : strain transformation
- constitutive equations : material model

Majority of engineering materials :

- homogenous : properties are the same in all locations
- isotropic : properties are the same in all directions
- linear : linear relationship between stress & strain
- elastic : no energy is lost in deformation



Cauchy Strain & Stress Tensors :

$$\bar{\sigma} = \begin{bmatrix} \sigma_{xx} & \tau_{xy} & \tau_{xz} \\ \tau_{xy} & \sigma_{yy} & \tau_{yz} \\ \tau_{xz} & \tau_{yz} & \sigma_{zz} \end{bmatrix} \quad \bar{\epsilon} = \begin{bmatrix} \epsilon_{xx} & \epsilon_{xy} & \epsilon_{xz} \\ \epsilon_{xy} & \epsilon_{yy} & \epsilon_{yz} \\ \epsilon_{xz} & \epsilon_{yz} & \epsilon_{zz} \end{bmatrix} \quad \text{note : } \epsilon_{xy} = \frac{\gamma_{xy}}{2}$$

- Linear material described by 6x6 matrix S with 36 material constants : S_{ij}

$$\begin{bmatrix} \epsilon_{xx} \\ \epsilon_{yy} \\ \epsilon_{zz} \\ \gamma_{yz} \\ \gamma_{xz} \\ \gamma_{xy} \end{bmatrix} = \begin{bmatrix} S_{11} & S_{12} & S_{13} & S_{14} & S_{15} & S_{16} \\ S_{21} & S_{22} & S_{23} & S_{24} & S_{25} & S_{26} \\ S_{31} & S_{32} & S_{33} & S_{34} & S_{35} & S_{36} \\ S_{41} & S_{42} & S_{43} & S_{44} & S_{45} & S_{46} \\ S_{51} & S_{52} & S_{53} & S_{54} & S_{55} & S_{56} \\ S_{61} & S_{62} & S_{63} & S_{64} & S_{65} & S_{66} \end{bmatrix} \begin{bmatrix} \sigma_{xx} \\ \sigma_{yy} \\ \sigma_{zz} \\ \tau_{yz} \\ \tau_{xz} \\ \tau_{xy} \end{bmatrix}$$

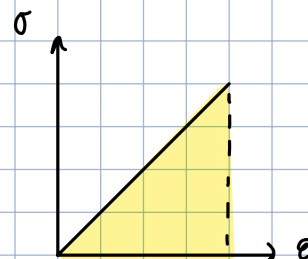
We want to reduce number of these constants.

Strain Energy: (Hooke's Law)

1. Apply stress σ_{xx} with resulting strain energy :

$$U_A = \frac{1}{2} \sigma_{xx} \epsilon_{xx} = \frac{1}{2} S_{11} \sigma_{xx}^2$$

↑
strain energy



strain energy = area under graph

2. Applying stress σ_{yy} while maintaining stress σ_{xx} :

$$U_b = \frac{1}{2} S_{22} \sigma_{yy}^2 + \underbrace{\sigma_{xx} S_{12} \sigma_{yy}}_{\substack{\epsilon_{xx} \text{ strain due} \\ \text{to } \sigma_{yy} \text{ (Poisson)}}} \leftarrow \text{units energy per unit volume}$$

Total elastic strain energy:

$$U_A + U_B = \frac{1}{2} S_{11} \sigma_{xx}^2 + \frac{1}{2} S_{22} \sigma_{yy}^2 + \sigma_{xx} S_{12} \sigma_{yy} = U$$

Repeat process applying σ_{yy} first then σ_{xx} :

$$U = \frac{1}{2} S_{22} \sigma_{yy}^2 + \frac{1}{2} S_{11} \sigma_{xx}^2 + \sigma_{xx} S_{12} \sigma_{yy}$$

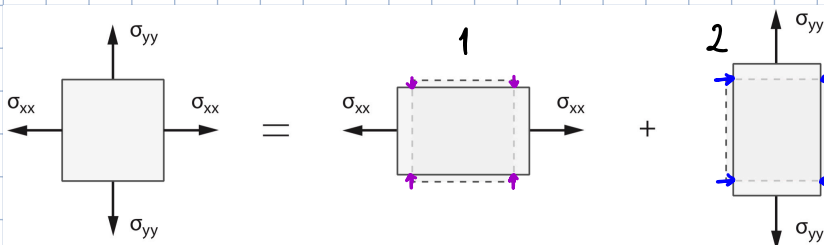
→ same result, ∴ **strain energy independent of order of application**

∴ $S_{ij} = S_{ji} \rightarrow 21$ elastic constants

Using mathematical approach to isotropic materials (too complex):

$$\begin{bmatrix} \epsilon_{xx} \\ \epsilon_{yy} \\ \epsilon_{zz} \\ \gamma_{yz} \\ \gamma_{xz} \\ \gamma_{xy} \end{bmatrix} = \begin{bmatrix} S_{11} & S_{12} & S_{12} & 0 & 0 & 0 \\ S_{12} & S_{11} & S_{12} & 0 & 0 & 0 \\ S_{12} & S_{12} & S_{11} & 0 & 0 & 0 \\ 0 & 0 & 0 & 2(S_{11} - S_{12}) & 0 & 0 \\ 0 & 0 & 0 & 0 & 2(S_{11} - S_{12}) & 0 \\ 0 & 0 & 0 & 0 & 0 & 2(S_{11} - S_{12}) \end{bmatrix} \begin{bmatrix} \sigma_{xx} \\ \sigma_{yy} \\ \sigma_{zz} \\ \tau_{yz} \\ \tau_{xz} \\ \tau_{xy} \end{bmatrix} \quad \text{2 constants: } S_{11} \text{ \& } S_{12}$$

Consider Bi-Axial Stress State:



Direct strains:

$$1. \quad \epsilon_{xx} = \frac{\sigma_{xx}}{E} \quad \& \quad \text{strain in y due to Poisson:} \quad \epsilon_{yy} = -\nu \epsilon_{xx} = -\nu \frac{\sigma_{xx}}{E}$$

$$2. \quad \epsilon_{yy} = \frac{\sigma_{yy}}{E} \quad \& \quad \text{strain in x due to Poisson:} \quad \epsilon_{xx} = -\nu \epsilon_{yy} = -\nu \frac{\sigma_{yy}}{E}$$

Superimposing two uni-axial strains:

$$\epsilon_{xx} = \frac{1}{E} (\sigma_{xx} - \nu \sigma_{yy})$$

$$\epsilon_{yy} = \frac{1}{E} (\sigma_{yy} - \nu \sigma_{xx})$$

→ $\epsilon_{xx} \neq -\nu \epsilon_{yy}$ for bi-axial stress state

↪ only applies to uni-axial

Shear Strains:

$$\gamma_{xy} = \frac{\tau_{xy}}{G}$$

→ combining into matrix formulation:

$$\begin{bmatrix} \epsilon_{xx} \\ \epsilon_{yy} \\ \gamma_{xy} \end{bmatrix} = \begin{bmatrix} 1/E & -\nu/E & 0 \\ -\nu/E & 1/E & 0 \\ 0 & 0 & 1/G \end{bmatrix} \begin{bmatrix} \sigma_{xx} \\ \sigma_{yy} \\ \tau_{xy} \end{bmatrix}$$

no link between applied shear stress and direct strain

no link between applied direct stress and shear strain

for isotropic materials, direct and shear strains decoupled

∴ principal directions for stress & strain coincide

Shear Modulus:

- for isotropic materials shear modulus G is not independent.

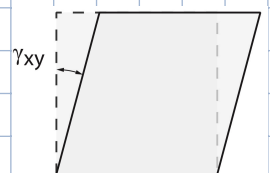
→ can see this if we consider that pure shear = bi-axial stress at 45° .

Shear strain due to pure shear:

$$\gamma_{xy} = \frac{\tau_{xy}}{G}$$

transforming to $X'Y'$ coordinate system ($+45^\circ$)

$$\epsilon_{x'x'} = \sin\theta \cos\theta \gamma_{xy} = \frac{\tau_{xy}}{2G} \quad (1)$$



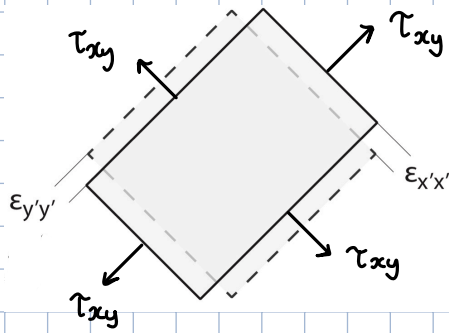
σ_{yy} & $\sigma_{xx} = 0$ as pure shear

Shear strain due to equivalent direct stresses:

$$\epsilon_{x'x'} = \frac{1}{E} (\sigma_{x'x'} - \nu \sigma_{y'y'})$$

$\sigma_{x'x'} = \sigma_{y'y'} = \tau_{xy}$ at 45°

$$= \frac{\tau_{xy}}{E} (1 - \nu) \quad (2)$$



equating both strains (1) & (2)

$$G = \frac{E}{2(1+\nu)}$$

for an isotropic, linear-elastic material under plane stress

Compliance matrix:

$$\begin{bmatrix} \epsilon_{xx} \\ \epsilon_{yy} \\ \gamma_{xy} \end{bmatrix} = \frac{1}{E} \begin{bmatrix} 1 & -\nu & 0 \\ -\nu & 1 & 0 \\ 0 & 0 & 2(1+\nu) \end{bmatrix} \begin{bmatrix} \sigma_{xx} \\ \sigma_{yy} \\ \tau_{xy} \end{bmatrix}$$

↓
Constitutive equations
↑

Stiffness matrix:

$$\begin{bmatrix} \sigma_{xx} \\ \sigma_{yy} \\ \tau_{xy} \end{bmatrix} = \frac{E}{1-\nu^2} \begin{bmatrix} 1 & \nu & 0 \\ \nu & 1 & 0 \\ 0 & 0 & (1-\nu)/2 \end{bmatrix} \begin{bmatrix} \epsilon_{xx} \\ \epsilon_{yy} \\ \gamma_{xy} \end{bmatrix}$$

↘ inverse

Plane Stress vs. Plane Strain:

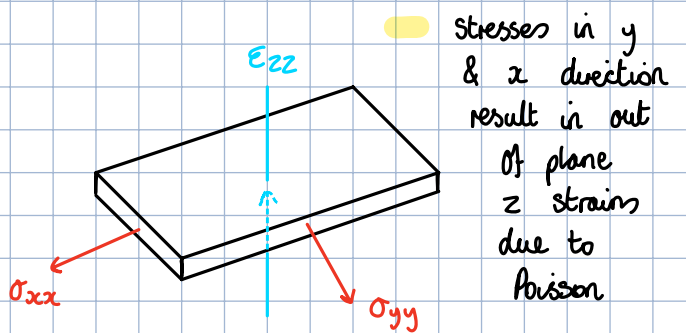
- expressions for through-thickness stress & strain (σ_{zz} & ϵ_{zz}):

Plane Stress: $\sigma_{zz} = \tau_{xz} = \tau_{yz} = 0$

→ through-thickness strain $\neq 0$

$$\epsilon_{zz} = \frac{1}{E} (\sigma_{zz} - \nu (\sigma_{xx} + \sigma_{yy}))$$

$$\epsilon_{zz} = -\frac{\nu}{E} (\sigma_{xx} + \sigma_{yy}) \neq 0$$



Plane Strain: $\epsilon_{zz} = \gamma_{xz} = \gamma_{yz} = 0$

$$\epsilon_{zz} = \frac{1}{E} (\sigma_{zz} - \nu(\sigma_{xx} + \sigma_{yy})) = 0$$

$$\rightarrow \sigma_{zz} = \nu(\sigma_{xx} + \sigma_{yy})$$

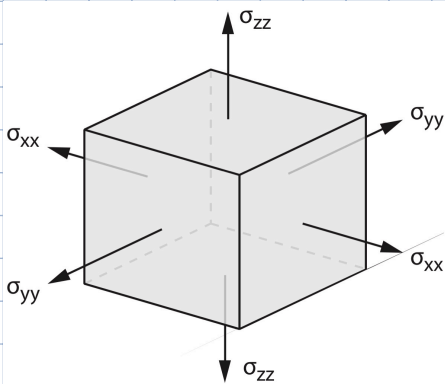
→ shows to ensure case of plane strain, you need through-thickness stress σ_{zz} .

Bulk Modulus:

- Only need two elastic constants to describe a material's properties:

↙
E & ν used in engineering

- by deriving bulk modulus, k , we can get further insight into material properties.



Volumetric strain describes change in material volume:

$$\frac{\Delta V}{V_0} = \frac{V - V_0}{V}$$

$$\text{where } V_0 = dx \, dy \, dz$$

Deformed Volume, V :

$$V = (1 + \epsilon_{xx}) dx (1 + \epsilon_{yy}) dy (1 + \epsilon_{zz}) dz = (1 + \epsilon_{xx})(1 + \epsilon_{yy})(1 + \epsilon_{zz}) dx \, dy \, dz$$
$$= (1 + \epsilon_{xx} + \epsilon_{yy} + \epsilon_{zz} + \text{higher order terms}) dx \, dy \, dz$$

$$\therefore \frac{\Delta V}{V_0} = \epsilon_{xx} + \epsilon_{yy} + \epsilon_{zz} = \frac{(\sigma_{xx} + \sigma_{yy} + \sigma_{zz})(1 - 2\nu)}{E}$$

Special case where tri-axial stress is spherical stress: $\sigma_{xx} = \sigma_{yy} = \sigma_{zz} = \sigma$

$$\frac{\Delta V}{V} = \frac{3\sigma(1 - 2\nu)}{E}$$

Bulk modulus relates spherical stress to volumetric strain:

$$\sigma = K \frac{\Delta V}{V_0}$$

→

$$K = \frac{E}{3(1-2\nu)}$$

- Looking at bulk & shear moduli: (both stiffnesses)

$$K = \frac{E}{3(1-2\nu)}$$

$$G = \frac{E}{2(1+\nu)}$$

- for realistic materials, shear & bulk moduli must be positive & finite

→ bounds Poisson's Ratio values $\nu \in \langle -1, 0.5 \rangle$

↳ most engineering materials
 $\nu \in \langle 0.2, 0.5 \rangle$

Poisson 'Extremes':

- $\nu = 0.5$ → infinite bulk modulus → incompressible e.g. rubber
- $\nu = 0$ e.g. cork → doesn't expand in normal direction to applied stresses
- $\nu < 0$ → auxetics